

The Crystal and Molecular Structure of 2-Acetyl-3-indazolinone, C₉H₈N₂O₂

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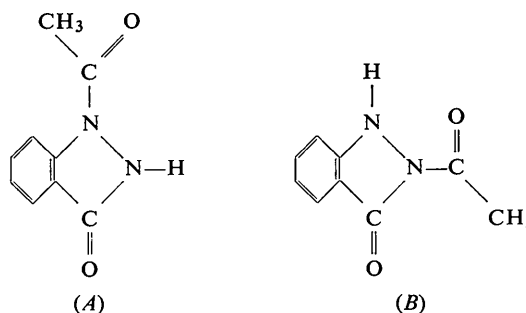
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The molecular structure of the unstable, low-melting (175–177°C) isomer of acetyl indazolinone, C₉H₈N₂O₂, has been determined from single-crystal, X-ray intensity data collected with an automated diffractometer. The crystals are monoclinic, space group $P2_1/n$, with four molecules in a unit cell of dimensions $a = 11.354$, $b = 5.683$, $c = 13.088$ Å; $\beta = 107.44^\circ$. The structure was solved by direct phasing of the structure amplitudes by the symbolic addition procedure and refined by full-matrix least-squares. Hydrogen atoms were located and refined. The final value of the discrepancy index is $R = 0.039$ for 1283 reflections. The results show that the unstable isomer is 2-acetyl-3-indazolinone.

Introduction

The use of 3-indazolinone and certain substituted 3-indazolinones as coupling agents for the production of color photographic images has been the subject of several patents (Barr, Henn & Salminen, 1962; Jennen & Eerdeken, 1954, 1960; Jennen, 1953, 1954). In particular, acetylation of 3-indazolinone under certain conditions leads to either of two isomers, 1-acetyl-3-indazolinone (*A*) or 2-acetyl-3-indazolinone (*B*) (Jennen, 1957; Jennen & Eerdeken, 1954). The lower melting isomer (hereafter called AIU) is unstable and is converted to the higher melting isomer (hereafter called AIS) in alkaline solution at room temperature or by heat (Young, 1965; Jennen, 1957). Reports in the literature (Janssen, 1962; Jennen, 1957) have supported the view that 2-acetyl-3-indazolinone (*B*) is the unstable isomer, AIU. However, some work in our laboratories (Young, 1965), involving direct synthesis of the couplers and the nature of the dyes formed on development, would be more easily explained if 2-acetyl-3-indazolinone (*B*) were the stable isomer, AIS.



To resolve this ambiguity in structure assignment, we have determined the crystal and molecular structure of AIU by single-crystal, X-ray diffraction techniques. The structure ultimately found for this isomer is 2-acetyl-3-indazolinone.

Experimental

Samples of both isomers were supplied by D. V. Young, of these laboratories. The sample of AIU already contained suitable single crystals; recrystallization of AIS

Table 1. *Crystal data for AIU*

C ₉ H ₈ N ₂ O ₂ , M.W. 176.18,	m.p. 175–177°
2-Acetyl-3-indazolinone.	
Colorless, irregular plates.	
Monoclinic (Cu $K\alpha_1 = 1.54051$ Å),	$a = 11.354 \pm 0.001$ Å
	$b = 5.683 \pm 0.001$
	$c = 13.088 \pm 0.001$
	$\beta = 107.44 \pm 0.01^\circ$
Volume of unit cell, 805.7 ± 0.3 Å ³ .	
Density, calculated ($Z = 4$)	1.452 g.cm ⁻³
measured (flotation)	1.44 g.cm ⁻³
Linear absorption coefficient, $\mu = 8.9$ cm ⁻¹ (Cu $K\alpha$)	
Total number of electrons per unit cell $F(000) = 368$	
Absent spectra: $h0l$ for $h + l$ odd, $0k0$ for k odd	
Space group $P2_1/n$ (C^5_2h)	
General positions: $\pm(x, y, z); \frac{1}{2} + x, \frac{1}{2} - y, \frac{1}{2} + z$.	

from acetone yielded satisfactory crystals. Small crystals of both materials were attached to thin glass rods with Canada balsam. Space group extinctions and preliminary lattice dimensions were obtained from precession photographs.

(A), AIS

The colorless acicular (*b*) crystals (m.p. 213–216°C) are monoclinic with cell parameters:

$$a = 17.02 \pm 0.03, b = 5.29 \pm 0.01, c = 9.02 \pm 0.02 \text{ \AA}, \beta = 100.4 \pm 0.1^\circ.$$

Systematic absences are *h*0*l* for *h* + *l* odd and 0*k*0 for *k* odd. These absences are consistent with space group *P*2₁/*n* (*C*_{2h}⁵). Four molecules per unit cell give a calculated density of 1.46 g.cm⁻³ in agreement with the observed density of 1.47 g.cm⁻³ obtained by flotation. No further work has been done with this isomer since the crystals are small and of generally poor quality.

(B), AIU

Some difficulty was experienced in finding untwinned crystals of AIU, but eventually an untwinned, irregular, platy crystal approximately 0.57 × 0.34 × 0.05 mm was chosen for intensity data collection on an automatic Picker four-circle goniostat. Narrow source, open-counter ω scans (Furnas, 1957) through several reflections showed the mosaicity to range from 0.15° to 0.25° for various directions in the crystal. Twenty-five reflections at moderately high Bragg angles (Cu Kα radiation) were accurately centered through very narrow vertical and horizontal slits at a take-off angle of 0.5°. These observations were used as input to the computer program *PICK2* (Ibers, 1966) which refined the cell and orientation parameters by the method of least-squares and generated the cards to control the automated goniostat. The refined cell parameters as well as other crystal data are given in Table 1.

Intensity data were collected with Ni-filtered Cu radiation at a take-off angle of 2.0°. A scintillation detector was used and the pulse-height analyzer was set for an approximately 90% window. All 1377 unique reflections for 2θ < 130° were measured by the θ–2θ scan technique (Furnas, 1957) at a 2θ scan rate of 1.0° min⁻¹. The scan range varied from 1.6° at low 2θ to 3.7° at high 2θ. Stationary crystal-stationary counter background counts were taken for 40 sec at each end of the scan. For count rates above 13,000 c.p.s. brass attenuators (~68% attenuation) were automatically inserted in the diffracted beam by the Picker system.

To check electronic and crystal stability during the period of data collection, the intensity of the 216 reflection was measured every 50th reflection. The small (~2%) systematic decrease observed in this standard intensity over the period of data collection was not considered to be significant. The standard deviation of 216, calculated from 31 observations, was 2.7 times what one would expect on the basis of counting statistics alone.

The reduction of the intensity data to structure amplitudes was accomplished with the computer program *DACOR* written by one of us (DLS). The background for a reflection was approximated by a straight line between the two measured background points. The intensities, corrected for background, were next corrected for Lorentz and polarization effects but not for absorption. Standard deviations σ(*I*) were based on counting statistics and were corrected to σ(*F*). Of the 1377 measured reflections, 110 were less than 2σ(*I*) and were considered as unobserved. The intensities of these unobserved reflections were then set equal to 2σ(*I*) and corrected to structure amplitudes, *F*_{Hm}.

Determination and refinement of the structure

The structure was solved by the direct phasing of the structure amplitudes by the symbolic addition procedure of Karle & Karle (1966). The magnitudes of the normalized structure factors, |*E*_H|, were computed from

$$E_{\mathbf{H}}^2 = F_{\mathbf{H}}^2 / \varepsilon \sum_{j=1}^N f_{j\mathbf{H}}^2$$

where ε = 2 when *H* is *h*0*l* or 0*k*0, ε = 1 otherwise, *N* is the number of atoms in the unit cell, *f*_{*jH*} is the atomic scattering factor for the *j*th atom, and the *F*_{*H*}² have been placed on an absolute scale and corrected for thermal motion by means of a least-squares fit of a Wilson (1942) plot.

The computer program *ESORT*, written by one of us (DLS), was used to calculate the statistical averages shown in Table 2 for the normalized structure factors. Comparison of these values with theoretical values (Hauptman & Karle, 1953) confirmed the centric space group. Unobserved reflections were included in the Wilson plot and in the statistical averages as *E*_{im}²/3 (Hamilton, 1955). The same program listed all triplet interactions *E*_H, *E*_K, and *E*_{H–K} for all |*E*| > 2.0 and also those interactions of |*E*_H| > 1.6 with |*E*_K| and |*E*_{H–K}| > 1.4 to facilitate the application of the Σ₂ relationship,

$$sE_{\mathbf{H}} \sim s \sum E_{\mathbf{K}} E_{\mathbf{H}-\mathbf{K}},$$

of Hauptman & Karle (1953) used in the symbolic addition procedure.

Table 2. Statistical averages of |*E*|

	Experimental	Theoretical	
		Centric	Acentric
⟨ <i>E</i> ⟩	0.811	0.798	0.886
⟨ <i>E</i> ² – 1 ⟩	0.995	0.968	0.736
⟨ <i>E</i> ² ⟩	1.034	1.000	1.000

Three linearly independent reflections with large *E* magnitudes and many interactions were assigned positive signs to specify the origin. Later phases *a*, *b* and *c* were assigned in sequence to three additional reflections in order to proceed with the phase determination. With these six assignments (Table 3) the phases of 218

of the 237 reflections with $|E| > 1.4$ were expressed in terms of +, -, or functions of a, b , and c . No phase indication was accepted unless the probability that it was correct was greater than 0.97, as calculated from Woolfson's (1954) formula,

$$P + (E_{\mathbf{H}}) = \frac{1}{2} + \frac{1}{2} \tanh(\sigma_3/\sigma_2^{3/2}) |E_{\mathbf{H}}| \frac{\sum_{\mathbf{K}} E_{\mathbf{K}} E_{\mathbf{H}-\mathbf{K}}}{K},$$

where $\sigma_n = \sum_{j=1}^N Z_j^n$. From interactions among the signs

it soon became evident that a was probably -. Later evidence indicated that both b and c were probably +. The three-dimensional E map (Karle, Hauptman, Karle & Wing, 1958), calculated with this most consistent set of phases, yielded all 13 heavy atoms for which the discrepancy index $R_1 = \sum ||F_o| - |F_c^*|| / \sum |F_o|$ was 0.27.

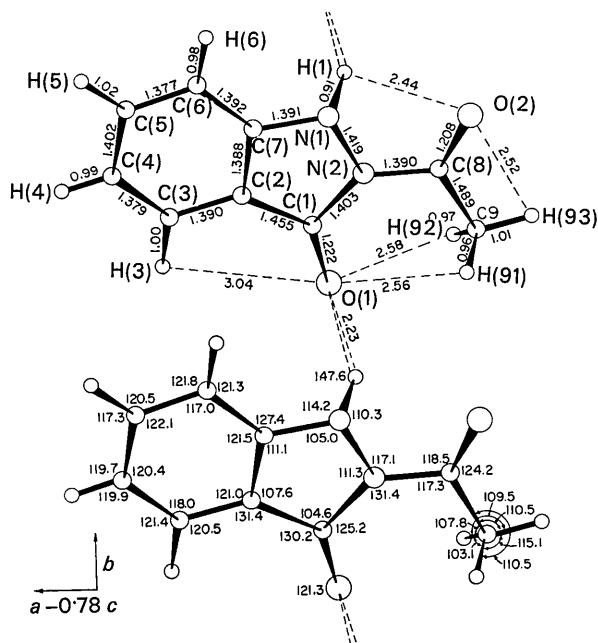


Fig. 1. Configuration, labeling, and interatomic distances and angles for AIU.

Table 3. Starting set of phases

h	k	l	$ E $	Phase
2	1	2	3.82	+
5	2	$\bar{1}$	2.98	+
8	1	$\bar{9}$	2.54	+
2	0	2	2.90	a
4	3	$\bar{1}$	3.09	b
2	1	$\bar{8}$	2.29	c

A full-matrix least-squares refinement was then made with the use of all observed reflections plus those unobserved reflections for which $|F_c^*| > F_{lim}$. The quantity minimized was $\sum w(|F_o| - |F_c^*|)^2$ where

$$|F_c^*| = KF_c \left[1 + g \left(\frac{1 + \cos^4 2\theta}{(1 + \cos^2 2\theta) \sin 2\theta} \right) F_c^2 \right]^{-1/2} = KF_c/G,$$

in which K is the scale factor, g is the extinction parameter (Zachariasen, 1963; Larson, 1967), and F_c is the

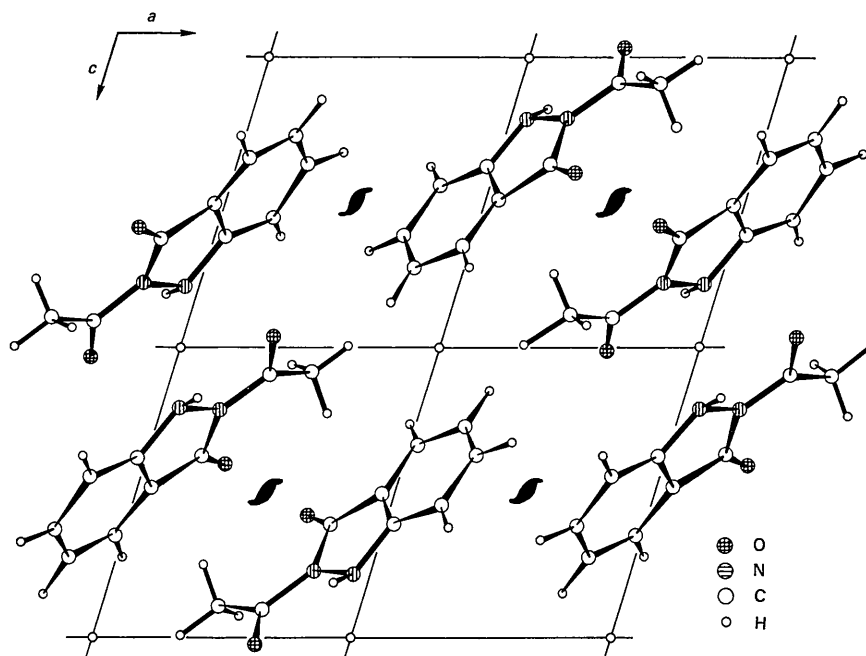


Fig. 2. $[010]_c$ projection of the structure.

usual calculated structure factor. The weights were defined as

$$w = [\sigma^2(F_o) + (0.03 F_o)^2]^{-1}.$$

Two cycles of least-squares refinement with isotropic temperature factors followed by two cycles with anisotropic temperature factors reduced R_1 to 0.080. A difference Fourier synthesis clearly showed the 8 hydrogen atoms. The hydrogen atoms were assigned isotropic temperature factors and refinement was resumed with all parameters except g allowed to vary. Four cycles resulted in $R_1 = 0.044$ with only small ($< 0.8\sigma$) parameter shifts on the last cycle. Up to this point, g had been held equal to zero. Two refinement cycles in which only K and g were varied reduced R_1 to 0.039 and two more cycles in which all parameters were varied concluded the refinement. For the last cycle of refinement only three of the 151 parameters shifted more than 0.1σ and the maximum shift was 0.29σ . The final discrepancy indices for the 1267 observed reflections plus the 16 unobserved reflections with $|F_c^*| > F_{lim}$ were $R_1 = 0.039$ and $R_2 = [\sum w(F_o - |F_c^*|)^2 / \sum w F_o^2]^{1/2} = 0.059$. The standard deviation of an observation of unit weight was 1.79. Of the 16 unobserved reflections with $|F_c^*| > F_{lim}$, none had $|F_c^*| > 1.5 F_{lim}$. It is interesting to note that all 218 phases originally derived from the Σ_2 relation were correct.

A final difference Fourier synthesis contained values between -0.20 and $+0.30 \text{ e.}\text{\AA}^{-3}$ and thus verified the structure. However, the difference electron density was generally > 0 in the vicinity of the mean molecular planes, particularly between atoms, and < 0 in the regions between the molecules. This systematic distribution of residual density is undoubtedly caused by the neglect of bonding effects on the atomic scattering factors; it

has been observed by Cady (1967), Fritchie (1966) and Hanson (1966), among others, and is discussed in detail by Hirshfeld & Rabinovich (1967), O'Connell, Rae & Maslen (1966), Rae & Maslen (1965), Dawson (1965) and Stewart, Davidson & Simpson (1965).

The final positional and thermal parameters, with standard deviations estimated from the full-matrix least-squares process, are given in Table 4. The average standard deviations of the positional parameters expressed in \AA are 0.0013 for O and N, 0.0015 for C, and 0.022 for H. Two other parameters refined in the least-squares process were the scale factor, $K = 8.013 \pm 0.038$, and the extinction coefficient,* $g = (0.269 \pm 0.031) + 10^{-4}$. The final values of the observed and calculated structure amplitudes are compared in Table 5.

The Fourier computer program used in this analysis was obtained from Professor J. Trotter (1965). The full-matrix least-squares refinement utilized a modification of the Busing, Martin & Levy (1962) program *ORFLS*. The atomic scattering factors used for all atoms were obtained from *International Tables for X-ray Crystallography* (1962). All computations were performed on an IBM System 360 computer.

Description and discussion of the structure

The structure determination has shown that AIU is 2-acetyl-3-indazolinone. Thus AIS, the stable isomer,

* A referee has noted that there has been an overcorrection for the extinction effect (see 202 and 212). This is undoubtedly a result of our method for determining g . We used a least-squares procedure in which all reflections were included and which was biased against the stronger reflections by the weighting scheme. Thus g was probably determined mainly by the reflections which were least affected by extinction.

Table 4. Final atomic parameters with their estimated standard deviations in parentheses

Thermal parameters are defined by the expression:

$$T_i = \exp[-\frac{1}{4}(B_{11}h^2a^*2 + B_{22}k^2b^*2 + B_{33}l^2c^*2 + 2B_{12}hka^*b^* + 2B_{13}hla^*c^* + 2B_{23}klb^*c^*)].$$

	x	y	z	B_{11}	B_{22}	B_{33}	B_{12}	B_{13}	B_{23}
O(1)	0.34275 (11)	-0.28796 (19)	0.79572 (9)	4.44 (5)	2.69 (5)	4.32 (5)	-0.72 (4)	1.45 (4)	-0.63 (4)
O(2)	0.31809 (13)	0.24484 (21)	1.01240 (10)	6.51 (8)	4.59 (6)	4.77 (6)	-1.23 (5)	3.25 (6)	-1.74 (5)
N(1)	0.46770 (11)	0.25925 (20)	0.89409 (10)	3.53 (6)	2.46 (5)	3.02 (5)	-0.51 (4)	1.00 (4)	-0.25 (4)
N(2)	0.38670 (10)	0.06660 (18)	0.88906 (9)	2.94 (5)	2.42 (5)	2.74 (5)	-0.26 (3)	0.76 (4)	-0.11 (3)
C(1)	0.39437 (12)	-0.09686 (23)	0.81098 (10)	2.90 (6)	2.44 (6)	2.72 (6)	0.21 (4)	0.45 (4)	-0.11 (4)
C(2)	0.47678 (12)	0.01186 (24)	0.75836 (10)	2.80 (5)	2.63 (6)	2.69 (5)	0.38 (4)	0.46 (4)	0.06 (4)
C(3)	0.51562 (14)	-0.06125 (28)	0.67234 (12)	3.69 (7)	3.38 (7)	3.31 (6)	0.60 (5)	1.03 (5)	-0.27 (5)
C(4)	0.59377 (14)	0.08456 (29)	0.63865 (13)	3.65 (7)	4.66 (8)	3.27 (6)	0.68 (5)	1.42 (5)	0.10 (5)
C(5)	0.63411 (14)	0.29745 (33)	0.69181 (12)	3.46 (6)	4.42 (8)	3.62 (7)	-0.01 (6)	1.19 (5)	0.64 (6)
C(6)	0.59648 (14)	0.37145 (27)	0.77726 (12)	3.52 (6)	3.31 (7)	3.30 (6)	-0.35 (5)	0.91 (5)	0.19 (5)
C(7)	0.51577 (12)	0.22491 (23)	0.80915 (11)	2.89 (6)	2.57 (6)	2.53 (5)	0.27 (4)	0.45 (4)	0.14 (4)
C(8)	0.31131 (13)	0.07519 (26)	0.95528 (11)	3.46 (6)	3.48 (7)	2.63 (6)	-0.01 (5)	0.83 (5)	-0.01 (5)
C(9)	0.22537 (16)	-0.12557 (33)	0.94959 (14)	3.70 (7)	4.42 (8)	3.72 (7)	-0.79 (6)	1.43 (6)	-0.34 (6)
H(1)	0.4267 (17)	0.3971 (32)	0.8914 (14)	1.88 (36)					
H(3)	0.4906 (17)	-0.2194 (33)	0.6386 (14)	1.64 (33)					
H(4)	0.6237 (19)	0.0365 (37)	0.5784 (16)	2.82 (40)					
H(5)	0.6956 (19)	0.3944 (34)	0.6666 (17)	2.83 (41)					
H(6)	0.6232 (18)	0.5227 (36)	0.8129 (16)	2.66 (38)					
H(91)	0.2701 (22)	-0.2715 (40)	0.9629 (17)	3.42 (48)					
H(92)	0.1748 (22)	-0.1411 (41)	0.8761 (19)	3.80 (49)					
H(93)	0.1708 (21)	-0.0951 (40)	0.9968 (19)	4.54 (55)					

Table 5 (cont.)

2-10	8* 5	4 9 50 53	6-10 [†]	48 46	9 -4 39 42	1 -8 25 41	4 1 8* 7	6 -6 13 7	1 5 28 27
2-11	29 29	4 -9 58 60	7 -11	10 9	9 -5 8 3	1 9 9 9	4 -1 23 23	6 -7 32 31	1 -5 8 1
2-12	40 38	4 -10 87 85	7 0 65 66	9 -6 37 37	1 -9 25 25	4 2 8 4	6 -8 15 12	1 6 72 71	
3 0	48 44	4 -11 13 12	7 1 55 57	9 -7 36 38	1 -10 63 61	4 -2 130 134	6 -9 26 25	1 -6 51 45	
3 1	67 67	4 -12 78 79	7 -1 34 36	9 -8 69 64	2 0 99 98	4 3 62 60	7 0 29 25	2 2 21 21	
3 -1	45 46	5 0 110 111	7 2 40 41	9 -9 25 24	2 1 131 131	4 -3 75 79	7 1 19 19	2 1 27 27	
3 2	66 64	5 1 80 80	7 -2 30 29	10 0 58 59	2 -1 31 30	4 4 87 87	7 -1 11 9	2 -1 8* 10	
3 -2	8* 3	5 -1 144 142	7 3 46 45	10 -1 8* 4	2 -2 119 117	4 -4 17 13	7 -2 13 12	2 -2 98 94	
3 3	129 131	5 2 47 46	7 -3 51 46	10 -2 58 58	2 -2 30 33	4 5 19 20	7 -2 46 48	2 -2 58 58	
3 3	89 86	5 -2 140 144	7 4 39 39	10 -3 45 42	2 3 37 39	4 -5 81 82	7 3 58 55	2 3 47 48	
3 4	197 191	5 3 71 75	7 -4 18 13	10 -4 67 68	2 -3 13 15	4 6 8 9	7 4 72 76	2 -3 33 35	
3 -4	117 119	5 -3 22 24	7 5 21 21	10 -5 21 22	2 4 58 59	4 -6 131 133	7 -4 39 39	2 4 79 77	
3 5	118 115	5 4 13 10	7 -5 40 47	10 -6 8* 1	2 4 66 64	4 7 32 32	7 -5 44 46	2 -4 88 91	
3 -5	31 33	5 -4 78 78	7 6 16 16	10 -7 22 18	2 5 81 83	4 -7 30 29	7 -6 29 24	2 -5 8* 2	
3 6	35 27	5 5 88 88	7 -6 136 134	2 -5 40 40	2 6 43 43	4 -8 8* 2	7 -7 43 42	2 -6 9 4	
3 -6	93 95	5 -5 128 128	7 -7 92 92	2 6 43 43	4 -9 58 59	8 0 17 16	7 -8 21 21	2 -6 9 4	
3 7	10 12	5 6 92 91	7 -8 82 84	2 7 12 10	4 -10 29 28	8 1 7* 1	8 0 43 42	3 -5 11 11	
3 -7	125 131	5 -6 23 23	7 -9 8* 3	2 7 12 10	5 0 88 90	8 1 7* 1	8 0 17 16	3 -5 11 11	
3 8	20 21	5 7 8* 5	7 -10 45 40	0 1 61 59	2 -7 75 75	5 1 10 10	8 -1 30 27	3 -1 16 16	
3 -8	85 94	5 -8 146 146	7 -11 10 9	1 2 125 127	3 3 84 84	5 6 29 25	8 -2 8 10	3 2 30 31	
3 9	45 48	5 8 14 13	8 0 10 13	0 3 23 22	2 -8 75 81	5 2 79 79	8 -3 54 54	3 -2 7* 3	
3 -9	73 71	5 -8 154 157	8 1 34 36	0 4 47 46	2 -9 28 27	5 -2 20 23	8 -4 8* 4	3 3 43 41	
3 10	19 17	5 -9 98 98	8 -1 29 32	0 5 21 22	2 -10 40 38	5 3 43 44	8 -5 53 51	3 3 8* 5	
3 -10	65 68	5 -10 72 72	8 2 60 63	0 6 36 40	3 0 157 151	5 -3 45 46	8 -6 8* 0	3 4 7* 9	
3 11	15 15	5 -11 49 45	8 -2 16 13	0 7 22 22	3 1 59 63	5 4 19 21		3 -4 42 44	
3 -12	14 13	5 -12 17 13	8 3 14 16	0 8 8* 3	3 1 32 31	5 4 76 77		3 -5 59 61	
4 0	141 138	6 0 210 212	8 -3 22 22	0 9 8* 7	3 2 22 22	5 5 11 13		3 -6 29 29	
4 1	72 70	6 1 74 72	8 4 17 17	1 0 63 62	3 -2 192 202	5 -5 88 90		4 0 25 23	
4 -1	71 73	6 -1 67 72	8 -4 43 42	1 1 149 150	3 3 84 86	5 6 29 25		4 1 42 41	
4 2	108 107	6 2 131 129	8 5 9 12	1 -1 77 78	3 -3 29 28	5 -6 73 73		4 -1 15 16	
4 -2	106 116	6 -2 22 20	8 -5 22 23	1 2 121 119	3 4 52 49	5 -7 17 19		4 2 23 20	
4 3	116 120	6 3 41 40	8 -6 79 79	1 -2 137 136	4 -4 147 149	5 -8 29 29		4 -2 44 45	
4 -3	13 8	6 -3 35 38	8 -7 52 49	1 3 136 129	3 7 27 25	5 -9 44 44		4 3 28 26	
4 4	63 64	6 4 21 18	8 -8 55 52	1 -3 42 39	3 -5 100 106	6 0 16 18		4 -3 30 30	
4 -4	44 44	6 -4 110 111	8 -9 74 72	1 4 42 42	3 6 8 7	6 1 42 42		4 -4 16 13	
4 5	36 40	6 5 102 101	8 -10 21 20	1 -4 127 127	3 7 9* 10	6 1 38 41		4 -5 5* 5	
4 -5	97 102	6 -5 26 25	9 0 14 16	1 5 28 26	3 8 22 24	6 2 53 55		4 -6 22 25	
4 6	6* 4	6 -6 37 35	9 1 8* 3	1 6 31 32	3 9 26 27	6 -2 19 10		4 -7 34 33	
4 -6	56 55	6 -7 36 35	9 -1 8* 3	1 6 31 32	3 9 26 27	6 -2 19 10		4 -8 47 47	
4 7	10 7	6 7 26 26	9 2 8 9	1 -6 159 156	3 -8 8* 2	6 -3 41 38		5 -1 9 11	
4 -7	48 48	6 -8 14 14	9 3 8* 8	1 7 50 49	3 -9 31 31	1 3 34 33		5 -2 54 54	
4 8	49 49	6 -8 14 14	9 3 8* 8	1 -7 89 92	3 -10 31 31	6 -4 4 5		5 -3 53 51	
4 -8	35 34	6 -9 50 47	9 -3 50 48	1 8 31 33	4 0 22 26	6 5 9 6		5 -4 21 20	
						6 -5 46 44		5 -5 18 15	

Table 6. Short intermolecular contacts

Hydrogen bond			
N(1)···O(1 ⁱⁱ)	= O(1)···N(1 ⁱⁱⁱ)	3·033 Å	
H(1)···O(1 ⁱⁱ)	= O(1)···H(1 ⁱⁱⁱ)	2·228	
Heavy atom contacts < 3·5 Å			
O(1)···N(2 ⁱⁱⁱ)	= N(2)···O(1 ^{iv})	3·089	
O(1)···C(1 ⁱⁱⁱ)	= C(1)···O(1 ^{iv})	3·167	
N(1)···C(8 ^v)	= C(8)···N(1 ^v)	3·291	
O(1)···C(8 ⁱⁱⁱ)	= C(8)···O(1 ^{iv})	3·318	
N(2)···N(2 ^v)		3·339	
O(1)···C(7 ⁱⁱⁱ)	= C(7)···O(1 ⁱⁱ)	3·369	
N(1)···N(2 ^v)	= N(2)···N(1 ^v)	3·377	
O(2)···C(6 ^{vi})	= C(6)···O(2 ^{vi})	3·417	
O(2)···C(1 ^v)	= C(1)···O(2 ^v)	3·491	
Heavy atom-hydrogen contacts < 3·0 Å			
O(2)···H(6 ^{vi})	= H(6)···O(2 ^{vi})	2·55	
C(7)···H(92 ^{iv})	= H(92)···C(7 ⁱⁱⁱ)	2·83	
O(2)···H(91 ⁱ)	= H(91)···O(2 ⁱⁱⁱ)	2·84	
C(2)···H(92 ^{iv})	= H(92)···C(2 ⁱⁱⁱ)	2·85	
O(2)···H(4 ^{vii})	= H(4)···O(2 ^{viii})	2·88	
O(2)···H(5 ^{vii})	= H(5)···O(2 ^{viii})	2·88	
C(6)···H(3 ⁱ)	= H(3)···C(6 ⁱⁱⁱ)	2·97	
O(1)···H(92 ⁱⁱⁱ)	= H(92)···O(1 ^{iv})	2·98	
Hydrogen-hydrogen contacts < 2·9 Å			
H(5)···H(93 ^{viii})	= H(93)···H(5 ^{vii})	2·44	
H(3)···H(6 ⁱⁱ)	= H(6)···H(3 ⁱ)	2·75	
H(4)···H(6 ^{ix})	= H(6)···H(4 ^x)	2·80	
H(4)···H(93 ^{xii})	= H(93)···H(4 ^{xii})	2·84	

i	x	1+y	z
ii	x	-1+y	z
iii	$\frac{1}{2}-x$	$-\frac{1}{2}+y$	$-\frac{1}{2}-z$
iv	$\frac{1}{2}-x$	$\frac{1}{2}+y$	$-\frac{1}{2}-z$
v	1-x	-y	2-z
vi	1-x	1-y	2-z
vii	$-\frac{1}{2}+x$	$\frac{1}{2}-y$	$\frac{1}{2}+z$
viii	$\frac{1}{2}+x$	$\frac{1}{2}-y$	$-\frac{1}{2}+z$
ix	$\frac{3}{2}-x$	$-\frac{1}{2}+y$	$\frac{3}{2}-z$
x	$\frac{3}{2}-x$	$\frac{1}{2}+y$	$\frac{3}{2}-z$
xi	$\frac{1}{2}+x$	$-\frac{1}{2}-y$	$-\frac{1}{2}+z$
xii	$-\frac{1}{2}+x$	$-\frac{1}{2}-y$	$\frac{1}{2}+z$

The shorter intermolecular contacts, listed in Table 6, are consistent with van der Waals radii (Pauling, 1960)

with the exception of H(1)···O(1ⁱⁱ) where a normal hydrogen bond occurs. The oxygen atom O(2) has several short intermolecular contacts but none short enough to be a hydrogen bond. Intramolecular hydrogen bonds which involve O(2) are ruled out since none of the intramolecular O···H distances is less than Bartell's (1960) hard-shell distances. Fig. 2 shows the packing of the molecules in the [010] projection of the structure. The hydrogen bonds join molecules in chains along *b*.

The bond distances and angles in AIU appear normal when compared with similar structures (Smith, 1969; Karle & Karle, 1967). The C-O bonds agree well with known ketonic C=O distances (Sutton, 1965). The fact that C(8)-O(2) is shorter than C(1)-O(1) is consistent with the hydrogen bonding. The C-H distances average 0·99 Å and N-H is 0·91 Å, both about 0·1 Å shorter than the true values but both agreeing with typical X-ray results. The C(8)-C(9) separation is normal for this type of single bond and the C(1)-C(2) bond (1·455 Å) is only slightly shorter than is considered normal (Lide, 1962). The N(1)-N(2) bond (1·419 Å) is somewhat shorter than the 1·45 Å for the N-N single bond in hydrazine (Sutton, 1965). The C-N bonds are all shorter than the unconjugated C-N single-bond length of 1·47 Å (Sutton, 1965) but are consistent with the observations of Karle & Karle (1967) for C-N bonds which involve trigonal carbon.

The C-C distances in the benzo ring average 1·388 Å and the C-C-C angles average 120·00°, both values in excellent agreement with the accepted values of 1·394 Å and 120·00° (Sutton, 1965). However, there is some systematic variation in these distances and angles. The bonds C(3)-C(4) and C(5)-C(6) (average 1·378 Å) are shorter than the bonds C(2)-C(3), C(6)-C(7) and C(7)-C(2) (average 1·390 Å), which are in turn shorter than C(4)-C(5) (1·402 Å). Likewise, the C-C-C angles at C(3) and C(6) are smaller than the other four angles of the benzo ring.

Table 7. Equations of least-squares planes and distances (Å) of atoms from these planes

X , Y and Z are expressed in Å units relative to the Cartesian axial system defined by the a , b and c^* axes. The planes were calculated by the method of Schomaker, Waser, Marsh & Bergman (1959) as modified by Blow (1960).

Plane A :	$-0.6160X + 0.4694Y - 0.6326Z = -7.4456$								
B :	$-0.6235X + 0.4636Y - 0.6296Z = -7.4476$								
C :	$-0.5115X + 0.4640Y - 0.7233Z = -8.3148$								
D :	$-0.6221X + 0.4646Y - 0.6302Z = -7.4473$								
E :	$-0.6038X + 0.4746Y - 0.6405Z = -7.5062$								
F :	$-0.5376X + 0.1410Y - 0.8313Z = -9.6608$								
Plane A			Plane B				Plane C		
N(1)	0.036	C(7)	-0.002	C(2)	0.003	*O(1)	0.047	O(2)	-0.001
N(2)	-0.044	*O(1)	0.083	C(3)	0.005	*C(1)	-0.009	N(1)	-0.002
C(1)	0.018	*C(8)	-0.231	C(4)	-0.009	*C(8)	-0.269	N(2)	0.001
C(2)	0.015	*H(1)	-0.633	C(5)	0.004	*H(1)	-0.646	C(8)	0.005
C(3)	0.010	*H(3)	0.073	C(6)	0.005	*H(3)	0.062	C(9)	-0.004
C(4)	-0.017	*H(4)	-0.010	C(7)	-0.007	*H(4)	0.004	*C(1)	-0.072
C(5)	-0.011	*H(5)	0.022	*N(1)	0.022	*H(5)	0.047	*C(7)	-0.229
C(6)	-0.003	*H(6)	-0.025	*N(2)	-0.071	*H(6)	-0.012	*H(1)	-0.622
Plane D			Plane E				Plane F		
C(1)	-0.003	*O(1)	0.054	N(1)	0.034	*O(1)	0.099	N(2)	0.000
C(2)	0.006	*N(1)	0.026	N(2)	-0.030	*C(3)	-0.022	C(7)	0.000
C(3)	0.006	*N(2)	-0.065	C(1)	0.024	*C(6)	-0.044	H(1)	0.000
C(4)	-0.010	*H(3)	0.064	C(2)	-0.001	*C(8)	-0.198	*N(1)	0.382
C(5)	0.002	*H(4)	0.002	C(7)	-0.022	*H(1)	-0.634		
C(6)	0.005	*H(5)	0.043						
C(7)	-0.005	*H(6)	-0.013						

* Not included in the calculation of the plane.

The overall pattern of bond lengths points to extensive conjugation of the molecule. Thus, it is not surprising to find that the atoms N(1), N(2), C(1)...C(7) which form the fused ring system all lie within 0.05 Å of a plane (A , Table 7). Consideration of the surrounding angles shows that C(1) and N(2) are coplanar with their immediate neighbors but that N(1) is tetrahedrally coordinated. In fact, N(1) lies 0.38 Å out of the plane (F , Table 7) of its neighbors, a value close to that for an ideal tetrahedron. Consequently, a better description of the molecule is given by two planes. The benzo group plus C(1) form a good plane (D , Table 7) and the group N(1), N(2), C(8), C(9), O(2) also form a good plane (C , Table 7). The dihedral angle between these planes is 8.3°.

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